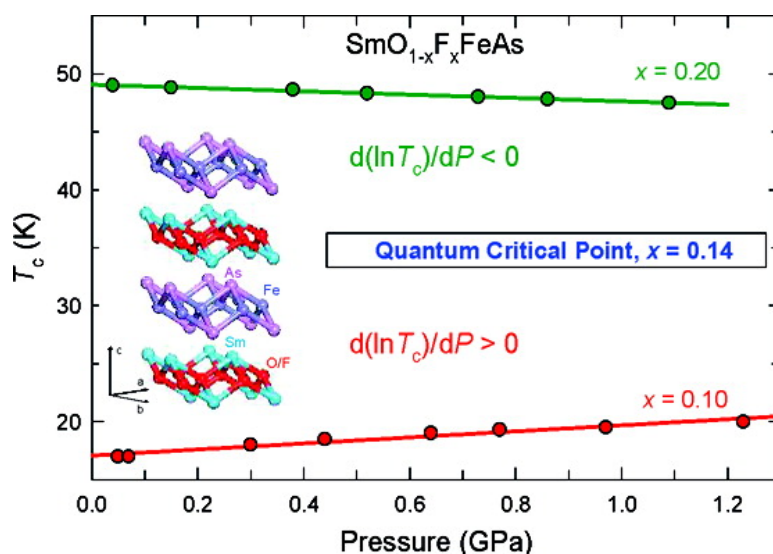


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Doping Dependence of the Pressure Response of T_c in the $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ Superconductors

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The discovery of superconductivity in the family of fluorine-doped rare-earth iron oxyarsenides, $\text{REO}_{1-x}\text{F}_x\text{FeAs}$ (RE = La, Ce, Pr, Nd, Sm, Gd) with transition temperatures, T_c as high as 55 K is currently generating considerable interest.^{1,2} These systems adopt tetragonal crystal structures (ZrCuSiAs-type, space group $P4/nmm$),¹ comprising interleaved two-dimensional RE(O,F) and FeAs layers (Figure 1 inset) with the latter providing the carrier conduction pathway. The appearance of high T_c superconductivity in these materials is generating exciting arguments on the possible mechanism of superconductivity and has led to considerable optimism for raising T_c to even higher values. In particular, it appears that superconductivity emerges at some critical value of F-doping when a competing spin-density-wave (SDW)-type long-range antiferromagnetic (AFM) order collapses; T_c increases smoothly with increasing doping level x until it reaches a maximum and then it decreases. This phenomenology together with the presence of magnetic interactions is reminiscent of the long established behavior of the high T_c cuprate superconductors and has led to suggestions of an unconventional non-BCS origin of the pairing mechanism.

A key experimental observation that can provide crucial information in differentiating between competing models of superconductivity in these oxyarsenides relates to the effect of applied pressure on T_c . Specifically it has been suggested that applied P should enhance the charge transfer between the insulating RE(O,F) and conducting FeAs slabs, thereby raising the value of T_c .³ Experimental data on the $\text{LaO}_{0.89}\text{F}_{0.11}\text{FeAs}$ superconductor appear to support this picture with T_c increasing with P at a rate of 1–2 K GPa^{-1} .^{3,4} Here we report an investigation of the magnetic response of the $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ ($0.05 \leq x \leq 0.20$) family of superconductors as a function of pressure up to 1.2 GPa. We find that at low F-doping levels ($x \leq 0.12$), the pressure coefficient, dT_c/dP is positive with T_c rapidly increasing with applied pressure. However, the sign of dT_c/dP sharply switches over and becomes negative as x increases beyond 0.15.

Polycrystalline samples with nominal composition $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ ($0 \leq x \leq 0.20$) were synthesized by conventional solid state reaction using high-purity SmAs, SmF_3 , Fe, and Fe_2O_3 , as described elsewhere.⁵ The samples were characterized by powder X-ray diffraction and temperature-dependent resistivity and dc magnetization measurements at ambient pressure. Bulk superconductivity is observed for $x = 0.10$ at ~ 12 K. T_c increases monotonically with increasing F content and reaches a maximum value of ~ 54 K (from resistivity measurements) at the optimal doping, $x = 0.20$ (Figure 1). This composition also represents the current upper limit of F doping in SmOFeAs .⁵ Magnetization measurements were carried

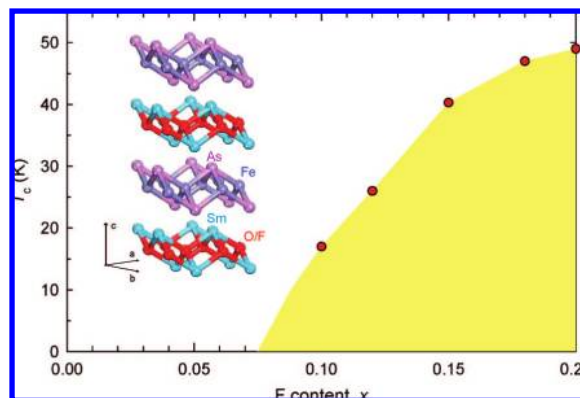


Figure 1. Doping dependence of the superconducting transition temperature T_c in $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$; (inset) schematic diagram of the tetragonal crystal structure of rare-earth iron oxyarsenides.

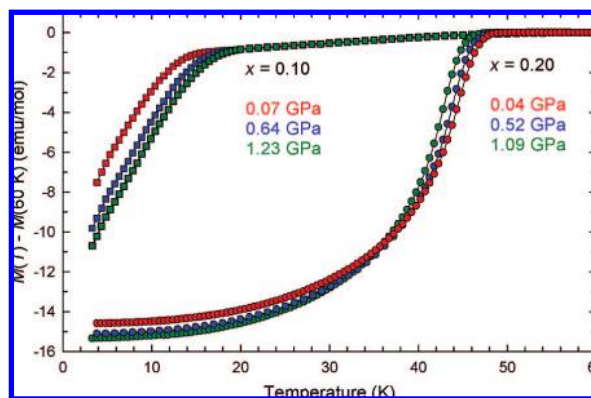


Figure 2. Temperature dependence of the magnetization, M (ZFC, 20 Oe) at selected pressures for $\text{SmO}_{0.90}\text{F}_{0.10}\text{FeAs}$ (squares) and $\text{SmO}_{0.80}\text{F}_{0.20}\text{FeAs}$ (circles).

out at 20 Oe on about 10-mg samples in the temperature range 1.8–60 K under both zero-field-cooling (ZFC) and field-cooling (FC) protocols with a Quantum Design SQUID magnetometer. Hydrostatic external pressure to ~ 1.2 GPa was applied with a piston–cylinder high-pressure cell (easyLab Technologies Mcell10) using high-purity Sn as an in situ manometer. The error in the pressure determination is on the order of 0.01–0.02 GPa. Daphne mineral oil was used as the pressure transmitting medium. Magnetization data were collected both on increasing and decreasing P .

Figure 2 shows the temperature dependence of the ZFC magnetization, M at various pressures for the $\text{SmO}_{0.90}\text{F}_{0.10}\text{FeAs}$ and $\text{SmO}_{0.80}\text{F}_{0.20}\text{FeAs}$ samples. Close to ambient pressure, bulk super-

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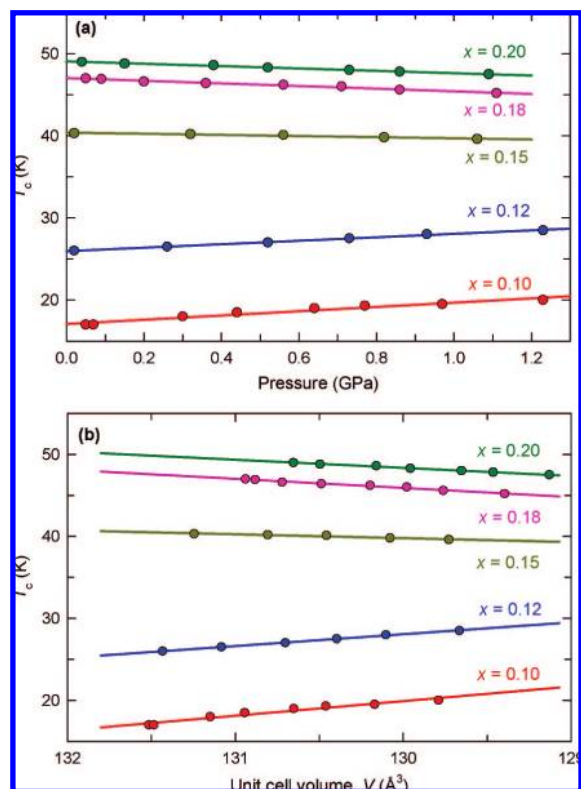


Figure 3. Superconducting transition temperature, T_c vs (a) applied pressure P and (b) unit-cell volume V for $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ ($0.10 \leq x \leq 0.20$). The lines through the points are linear fits to the data.

conductivity with onset T_{CS} of 17 and 49 K and shielding fractions of 12% and 23% (at 3.8 K) is observed for $x = 0.10$ and 0.20, respectively (Figure 2). However, the superconducting response of the samples to increasing external hydrostatic pressure is found to be drastically different. Although in both cases, T_c responds to pressure very sensitively, it shifts quasi-linearly to higher temperatures (at a rate of 2.6(2) K GPa^{-1}) for $\text{SmO}_{0.90}\text{F}_{0.10}\text{FeAs}$ and to lower temperatures (at a rate of $-1.44(5)$ K GPa^{-1}) for $\text{SmO}_{0.80}\text{F}_{0.20}\text{FeAs}$. Unlike earlier measurements,^{3,4,6} no broadening of the transitions is observed with increasing P implying the absence of sample inhomogeneities. The relation between T_c and pressure for $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ ($0.10 \leq x \leq 0.20$) is summarized in Figure 3a, which clearly shows that the change in sign of the pressure coefficient, dT_c/dP occurs for compositions with F doping levels between 0.12 and 0.15. We also note that low-field magnetization measurements on a $\text{SmO}_{0.95}\text{F}_{0.05}\text{FeAs}$ sample revealed the presence of only pressure-independent trace ($\sim 0.1\%$) superconductivity below 5 K and no pressure-induced transition to a superconducting state up to 1.2 GPa.

Of paramount importance is the determination of the underlying mechanism for superconductivity in the $\text{REO}_{1-x}\text{F}_x\text{FeAs}$ phases, and experimental data are rapidly being accumulated. Specifically for the $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ series, resistivity and Hall effect measurements⁵ have been interpreted in terms of the occurrence near $x \approx 0.14$ of a quantum critical point (QCP) owing to the competition of SDW order and superconductivity. Then it is indeed remarkable that the change in sign of dT_c/dP from positive to negative is observed in our experiments precisely at doping levels that straddle the putative QCP. Our experimental results are therefore not in agreement with theoretical models⁷ which have proposed that dT_c/dP should change sign from positive to negative when the compounds cross from the

so-called underdoped (T_c increases with x) to the overdoped (T_c decreases with x) region.

The absolute values of the normalized pressure coefficient, $d(\ln T_c)/dP$ also vary significantly across the $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ series. Close to the critical composition at which the superconducting state emerges, $d(\ln T_c)/dP$ assumes a maximum value of 0.15(1) GPa^{-1} . As the doping level increases further, $d(\ln T_c)/dP$ reaches a value of 0.081(2) GPa^{-1} at $x = 0.12$. Following the reversal in sign of the pressure coefficient, $d(\ln T_c)/dP$ is much smaller in magnitude at $-0.017(3)$ GPa^{-1} for $x = 0.15$; as x increases further and approaches optimal doping, $d(\ln T_c)/dP$ remains essentially constant at -0.03 GPa^{-1} . Of particular importance will be also to establish the doping dependence of the volume coefficients of T_c . At present, there are no experimental data available for the volume compressibility of $\text{REO}_{1-x}\text{F}_x\text{FeAs}$ phases but theoretical calculations⁸ provide a value of the bulk modulus B on the order of 90 GPa. Taking this into account, we estimate $ld(\ln T_c)/dV$ to be 0.104(8) \AA^{-3} at $x = 0.10$, implying a very sensitive dependence of the superconducting properties of $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ to the interatomic distances in the underdoped regime. The results of converting the $T_c(P)$ data in Figure 3a to $T_c(V)$ are shown in Figure 3b. As the doping level exceeds that of the QCP, $ld(\ln T_c)/dV$ is strongly suppressed and the dependence of the superconductivity onset to the interatomic distances is almost an order of magnitude smaller (0.012(2) \AA^{-3} at $x = 0.15$). By the time optimal doping is reached the volume coefficient of T_c is somewhat larger but still about five times smaller ($ld(\ln T_c)/dV = 0.020(1)$ \AA^{-3} at $x = 0.20$) than for that in the underdoped regime.

In conclusion, the observed doping dependence of the pressure coefficients in the $\text{SmO}_{1-x}\text{F}_x\text{FeAs}$ ($0.10 \leq x \leq 0.20$) superconductors reveals a sharp change in sign from positive to negative in the vicinity of the QCP proposed by resistivity and Hall effect measurements.⁵ In the underdoped region, the compounds display an extremely sensitive dependence of T_c on the interatomic distances. However, beyond the QCP and as optimal doping is approached, this sensitivity is strongly suppressed by a factor of 5–10. These results should form a stringent test of competing models for the interpretation of the superconducting pairing mechanism in fluorine-doped rare-earth iron oxyarsenides.

Supporting Information Available: Pressure-dependent magnetization measurements of the $x = 0.12$, 0.15, and 0.18 samples. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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